

Note:

CHAPTER 1: FLUIDINTRO

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This chapter is part of the textbook:

“Basics of Fluid Mechanics”

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MARCH 2, 2010

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NSY = Not Started Yet

CHAPTER 1

Introduction

1.1 What is Fluid Mechanics?

Fluid mechanics deals with the study of all fluids under static and dynamic situations. Fluid mechanics is a branch of continuous mechanics which deals with a relationship between forces, motions, and statical conditions in continuous material. This study area deals with many and diversified problems such as surface tension, fluid statics, flow in enclosed bodies, or flow round bodies (solid or otherwise), flow stability, etc. In fact, almost any action a person is doing involves some kind of a fluid mechanics problem. Furthermore, the boundary between the solid mechanics and fluid mechanics is some kind of gray shed and not a sharp distinction (see Figure 1.1 for the complex relationships between the different branches which only part of it should be drawn in the same time.). For example, glass appears as a solid material, but a closer look reveals that the glass is a liquid with a large viscosity. A proof of the glass “liquidity” is the change of the glass thickness in high windows in European Churches after hundred years. The bottom part of the glass is thicker than the top part. Materials like sand (some call it quick sand) and grains should be treated as liquids. It is known that these materials have the ability to drown people. Even material such as aluminum just below the mushy zone also behaves as a liquid similarly to butter. After it was established that the boundaries of fluid mechanics aren’t sharp, the discussion in this book is limited to simple and (mostly) Newtonian (sometimes power fluids) fluids which will be defined later.

The fluid mechanics study involve many fields that have no clear boundary between them. Researchers distinguish between orderly flow and chaotic flow as the laminar flow and the turbulent flow. The fluid mechanics can also be distinguish between a single phase flow and multiphase flow (flow made more than one phase or single distinguishable material). The last boundary (as all the boundaries in fluid mechanics)

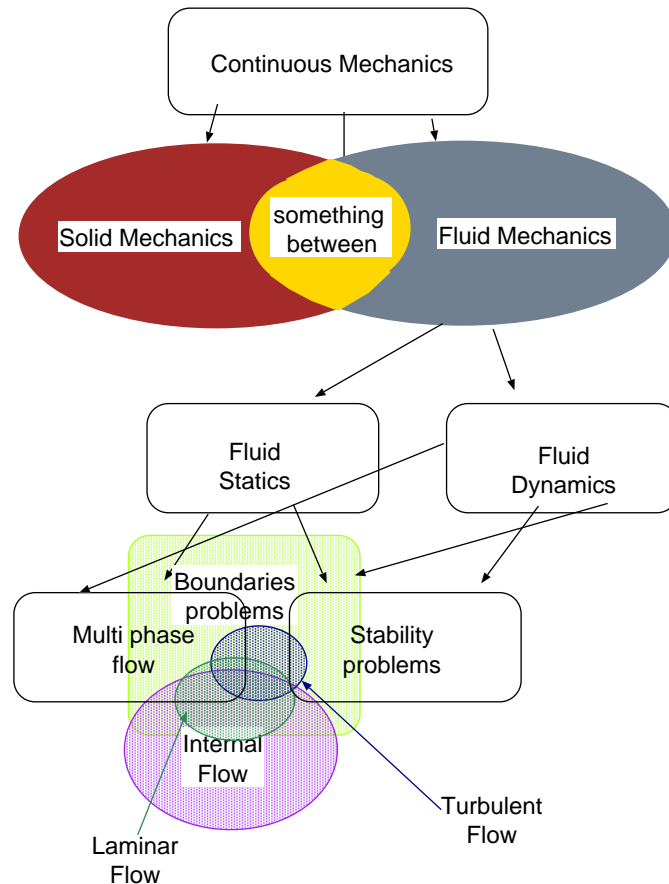


Fig. -1.1. Diagram to explain part of relationships of fluid mechanics branches.

isn't sharp because fluid can go through a phase change (condensation or evaporation) in the middle or during the flow and switch from a single phase flow to a multi phase flow. Moreover, flow with two phases (or materials) can be treated as a single phase (for example, air with dust particle).

After it was made clear that the boundaries of fluid mechanics aren't sharp, the study must make arbitrary boundaries between fields. Then the dimensional analysis will be used explain why in certain cases one distinguish area/principle is more relevant than the other and some effects can be neglected. Or, when a general model is need because more parameters are effecting the situation. It is this author's personal experience that the knowledge and ability to know in what area the situation lay is one of the main problems. For example, engineers in software company (EKK Inc, <http://ekkin.com/HTML>) analyzed a flow of a complete still liquid assuming a

complex turbulent flow model. Such absurd analysis are common among engineers who do not know which model can be applied. Thus, one of the main goals of this book is to explain what model should be applied. Before dealing with the boundaries, the simplified private cases must be explained.

There are two main approaches of presenting an introduction of fluid mechanics materials. The first approach introduces the fluid kinematic and then the basic governing equations, to be followed by stability, turbulence, boundary layer and internal and external flow . The second approach deals with the Integral Analysis to be followed with Differential Analysis, and continue with Empirical Analysis. These two approaches pose a dilemma to anyone who writes an introductory book for the fluid mechanics. These two approaches have justifications and positive points. Reviewing many books on fluid mechanics made it clear, there isn't a clear winner. This book attempts to find a hybrid approach in which the kinematic is presented first (aside to standard initial four chapters) follow by Integral analysis and continued by Differential analysis. The ideal flow (frictionless flow) should be expanded compared to the regular treatment. This book is unique in providing chapter on multiphase flow. Naturally, chapters on open channel flow (as a sub class of the multiphase flow) and compressible flow (with the latest developments) are provided.

1.2 *Brief History*

The need to have some understanding of fluid mechanics started with the need to obtain water supply. For example, people realized that wells have to be dug and crude pumping devices need to be constructed. Later, a large population created a need to solve waste (sewage) and some basic understanding was created. At some point, people realized that water can be used to move things and provide power. When cities increased to a larger size, aqueducts were constructed. These aqueducts reached their greatest size and grandeur in those of the City of Rome and China.

Yet, almost all knowledge of the ancients can be summarized as application of instincts, with the exception Archimedes (250 B.C.) on the principles of buoyancy. For example, larger tunnels built for a larger water supply, etc. There were no calculations even with the great need for water supply and transportation. The first progress in fluid mechanics was made by Leonardo Da Vinci (1452-1519) who built the first chambered canal lock near Milan. He also made several attempts to study the flight (birds) and developed some concepts on the origin of the forces. After his initial work, the knowledge of fluid mechanics (hydraulic) increasingly gained speed by the contributions of Galileo, Torricelli, Euler, Newton, Bernoulli family, and D'Alembert. At that stage theory and experiments had some discrepancy. This fact was acknowledged by D'Alembert who stated that, "The theory of fluids must necessarily be based upon experiment." For example the concept of ideal liquid that leads to motion with no resistance, conflicts with the reality.

This discrepancy between theory and practice is called the "D'Alembert paradox" and serves to demonstrate the limitations of theory alone in solving fluid problems. As in thermodynamics, two different of school of thoughts were created: the first be-

lied that the solution will come from theoretical aspect alone, and the second believed that solution is the pure practical (experimental) aspect of fluid mechanics. On the theoretical side, considerable contribution were made by Euler, La Grange, Helmholtz, Kirchhoff, Rayleigh, Rankine, and Kelvin. On the "experimental" side, mainly in pipes and open channels area, were Brahm's, Bossut, Chezy, Dubuat, Fabre, Coulomb, Dupuit, d'Aubisson, Hagen, and Poisseuille.

In the middle of the nineteenth century, first Navier in the molecular level and later Stokes from continuous point of view succeeded in creating governing equations for real fluid motion. Thus, creating a matching between the two school of thoughts: experimental and theoretical. But, as in thermodynamics, people cannot relinquish control. As results it created today "strange" names: Hydrodynamics, Hydraulics, Gas Dynamics, and Aeronautics.

The Navier-Stokes equations, which describes the flow (or even Euler equations), were considered unsolvable during the mid nineteenth century because of the high complexity. This problem led to two consequences. Theoreticians tried to simplify the equations and arrive at approximated solutions representing specific cases. Examples of such work are Hermann von Helmholtz's concept of vortexes (1858), Lanchester's concept of circulatory flow (1894), and the Kutta-Joukowski circulation theory of lift (1906). The experimentalists, at the same time proposed many correlations to many fluid mechanics problems, for example, resistance by Darcy, Weisbach, Fanning, Ganguillet, and Manning. The obvious happened without theoretical guidance, the empirical formulas generated by fitting curves to experimental data (even sometime merely presenting the results in tabular form) resulting in formulas that the relationship between the physics and properties made very little sense.

At the end of the twenty century, the demand for vigorous scientific knowledge that can be applied to various liquids as opposed to formula for every fluid was created by the expansion of many industries. This demand coupled with new several novel concepts like the theoretical and experimental researches of Reynolds, the development of dimensional analysis by Rayleigh, and Froude's idea of the use of models change the science of the fluid mechanics. Perhaps the most radical concept that effects the fluid mechanics is of Prandtl's idea of boundary layer which is a combination of the modeling and dimensional analysis that leads to modern fluid mechanics. Therefore, many call Prandtl as the father of modern fluid mechanics. This concept leads to mathematical basis for many approximations. Thus, Prandtl and his students Blasius, von Karman, Meyer, and Blasius and several other individuals as Nikuradse, Rose, Taylor, Bhuckingham, Stanton, and many others, transformed the fluid mechanics to modern science that we have known today.

While the understanding of the fundamentals did not change much, after World War Two, the way how it was calculated changed. The introduction of the computers during the 60s and much more powerful personal computer has changed the field. There are many open source programs that can analyze many fluid mechanics situations. Today many problems can be analyzed by using the numerical tools and provide reasonable results. These programs in many cases can capture all the appropriate parameters and adequately provide a reasonable description of the physics. However, there are many

other cases that numerical analysis cannot provide any meaningful result (trends). For example, no weather prediction program can produce good engineering quality results (where the snow will fall within 50 kilometers accuracy. Building a car with this accuracy is a disaster). In the best scenario, these programs are as good as the input provided. Thus, assuming turbulent flow for still flow simply provides erroneous results (see for example, EKK, Inc).

1.3 Kinds of Fluids

Some differentiate fluid from solid by the reaction to shear stress. It is a known fact said that the fluid continuously and permanently deformed under shear stress while solid exhibits a finite deformation which does not change with time. It is also said that liquid cannot return to their original state after the deformation. This differentiation leads to three groups of materials: solids and liquids. This test creates a new material group that shows dual behaviors; under certain limits; it behaves like solid and under others it behaves like liquid (see Figure 1.1). The study of this kind of material called rheology and it will (almost) not be discussed in this book. It is evident from this discussion that when a liquid is at rest, no shear stress is applied.

The fluid is mainly divided into two categories: liquids and gases. The main difference between the liquids and gases state is that gas will occupy the whole volume while liquids has an almost fix volume. This difference can be, for most practical purposes considered, sharp even though in reality this difference isn't sharp. The difference between a gas phase to a liquid phase above the critical point are practically minor. But below the critical point, the change of water pressure by 1000% only change the volume by less than 1 percent. For example, a change in the volume by more 5% will required tens of thousands percent change of the pressure. So, if the change of pressure is significantly less than that, then the change of volume is at best 5%. Hence, the pressure will not affect the volume. In gaseous phase, any change in pressure directly affects the volume. The gas fills the volume and liquid cannot. Gas has no free interface/surface (since it does fill the entire volume).

There are several quantities that have to be addressed in this discussion. The first is **force** which was reviewed in physics. The unit used to measure is [N]. It must be remember that force is a vector, e.g it has a direction. The second quantity discussed here is the area. This quantity was discussed in physics class but here it has an additional meaning, and it is referred to the direction of the area. The direction of area is perpendicular to the area. The area is measured in [m^2]. Area of three-dimensional object has no single direction. Thus, these kinds of areas should be addressed infinitesimally and locally.

The traditional quantity, which is force per area has a new meaning. This is a result of division of a vector by a vector and it is referred to as tensor. In this book, the emphasis is on the physics, so at this stage the tensor will have to be broken into its components. Later, the discussion on the mathematic meaning will be presented (later version). For the discussion here, the pressure has three components, one in the area direction and two perpendicular to the area. The pressure component in the area

direction is called pressure (great way to confuse, isn't it?). The other two components are referred as the shear stresses. The units used for the pressure components is $[N/m^2]$.

The density is a property which requires that liquid to be continuous. The density can be changed and it is a function of time and space (location) but must have a continuous property. It doesn't mean that a sharp and abrupt change in the density cannot occur. It referred to density that is independent of the sampling size. Figure 1.2 shows the density as a function of the sample size. After certain sample size, the density remains constant. Thus, the density is defined as

$$\rho = \lim_{\Delta V \rightarrow \varepsilon} \frac{\Delta m}{\Delta V} \quad (1.1)$$

It must be noted that ε is chosen so that the continuous assumption is not broken, that is, it did not reach/reduced to the size where the atoms or molecular statistical calculations are significant (see Figure 1.2 for point where the green lines converge to constant density). When this assumption is broken, then, the principles of statistical mechanics must be utilized.

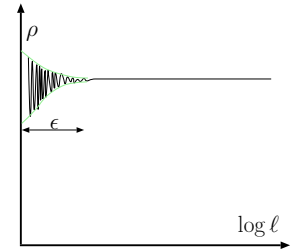


Fig. -1.2. Density as a function of the size of sample.

1.4 Shear Stress

The shear stress is part of the pressure tensor. However, here it will be treated as a separate issue. In solid mechanics, the shear stress is considered as the ratio of the force acting on area in the direction of the forces perpendicular to area. Different from solid, fluid cannot pull directly but through a solid surface. Consider liquid that undergoes a shear stress between a short distance of two plates as shown in Figure (1.3).

Fig. -1.3. Schematics to describe the shear stress in fluid mechanics.

The upper plate velocity generally will be

$$U = f(A, F, h) \quad (1.2)$$

Where A is the area, the F denotes the force, h is the distance between the plates. From solid mechanics study, it was shown that when the force per area increases, the velocity of the plate increases also. Experiments show that the increase of height will increase the velocity up to a certain range. Consider moving the plate with a zero lubricant ($h \sim 0$) (results in large force) or a large amount of lubricant (smaller force). In this discussion, the aim is to develop differential equation, thus the small distance analysis is applicable.

For cases where the dependency is linear, the following can be written

$$U \propto \frac{hF}{A} \tag{1.3}$$

Equations (1.3) can be rearranged to be

$$\frac{U}{h} \propto \frac{F}{A} \tag{1.4}$$

Shear stress was defined as

$$\tau_{xy} = \frac{F}{A} \tag{1.5}$$

From equations (1.4) and (1.5) it follows that ratio of the velocity to height is proportional to shear stress. Hence, applying the coefficient to obtain a new equality as

$$\tau_{xy} = \mu \frac{U}{h} \tag{1.6}$$

Where μ is called the absolute viscosity or dynamic viscosity which will be discussed later in this chapter in great length.

In steady state, the distance the upper plate moves after small amount of time, δt is

$$d\ell = U \delta t \tag{1.7}$$

From Figure 1.4 it can be noticed that for a small angle, the regular approximation provides

$$d\ell = U \delta t = \widehat{h \delta\beta} \tag{1.8}$$

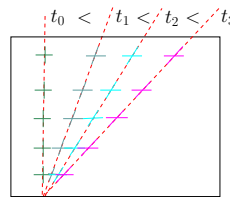


Fig. -1.4. The deformation of fluid due to shear stress as progression of time.

From equation (1.8) it follows that

$$U = h \frac{\delta\beta}{\delta t} \tag{1.9}$$

Combining equation (1.9) with equation (1.6) yields

$$\tau_{xy} = \mu \frac{\delta\beta}{\delta t} \tag{1.10}$$

If the velocity profile is linear between the plate (it will be shown later that it is consistent with derivations of velocity), then it can be written for small a angel that

$$\frac{\delta\beta}{\delta t} = \frac{dU}{dy} \tag{1.11}$$

Materials which obey equation (1.10) referred to as Newtonian fluid. For this kind of substance

$$\tau_{xy} = \mu \frac{dU}{dy} \quad (1.12)$$

Newtonian fluids are fluids which the ratio is constant. Many fluids fall into this category such as air, water etc. This approximation is appropriate for many other fluids but only within some ranges.

Equation (1.9) can be interpreted as momentum in the x direction transferred into the y direction. Thus, the viscosity is the resistance to the flow (flux) or the movement. The property of viscosity, which is exhibited by all fluids, is due to the existence of cohesion and interaction between fluid molecules. These cohesion and interactions hamper the flux in y -direction. Some referred to shear stress as viscous flux of x -momentum in the y -direction. The units of shear stress are the same as flux per time as following

$$\frac{F}{A} \left[\frac{kg \ m}{sec^2 \ m^2} \right] = \frac{\dot{m}U}{A} \left[\frac{kg \ m}{sec \ sec \ m^2} \right]$$

Thus, the notation of τ_{xy} is easier to understand and visualize. In fact, this interpretation is more suitable to explain the molecular mechanism of the viscosity. The units of absolute viscosity are $[N \ sec/m^2]$.

Example 1.1:

A space of 1 [cm] width between two large plane surfaces is filled with glycerine. Calculate the force that is required to drag a very thin plate of 1 [m²] at a speed of 0.5 m/sec. It can be assumed that the plates remains in equidistance from each other and steady state is achieved instantly.

SOLUTION

Assuming Newtonian flow, the following can be written (see equation (1.6))

$$F = \frac{A\mu U}{h} \sim \frac{1 \times 1.069 \times 0.5}{0.01} = 53.45[N]$$

End Solution

Example 1.2:

Castor oil at 25°C fills the space between two concentric cylinders of 0.2[m] and 0.1[m] diameters with height of 0.1 [m]. Calculate the torque required to rotate the inner cylinder at 12 rpm, when the outer cylinder remains stationary. Assume steady state conditions.

SOLUTION

The velocity is

$$U = r \dot{\theta} = 2 \pi r_i \text{ rps} = 2 \times \pi \times 0.1 \times \overbrace{12/60}^{\text{rps}} = 0.4 \pi r_i$$

Where *rps* is revolution per second.

The same way as in example (1.1), the moment can be calculated as the force times the distance as

$$M = F \ell = \frac{\overbrace{r_i} \overbrace{2 \pi r_i h} A \mu U}{r_o - r_i}$$

In this case $r_o - r_i = h$ thus,

$$M = \frac{2 \pi^2 \overbrace{0.1^3}^{r_i} \cancel{h} \overbrace{0.986}^{\mu} 0.4}{\cancel{h}} \sim .0078 [N m]$$

End Solution

1.5 Viscosity

1.5.1 General

Viscosity varies widely with temperature. However, temperature variation has an opposite effect on the viscosities of liquids and gases. The difference is due to their fundamentally different mechanism creating viscosity characteristics. In gases, molecules are sparse and cohesion is negligible, while in the liquids, the molecules are more compact and cohesion is more dominant. Thus, in gases, the exchange of momentum between layers brought as a result of molecular movement normal to the general direction of flow, and it resists the flow.

This molecular activity is known to increase with temperature, thus, the viscosity of gases will increase with temperature. This reasoning is a result of the considerations of the kinetic theory. This theory indicates that gas viscosities vary directly with the square root of temperature. In liquids, the momentum exchange due to molecular movement is small compared to the cohesive forces between the molecules. Thus, the viscosity is primarily dependent on the magnitude of these cohesive forces. Since these forces decrease rapidly with increases of temperature, liquid viscosities decrease as temperature increases.

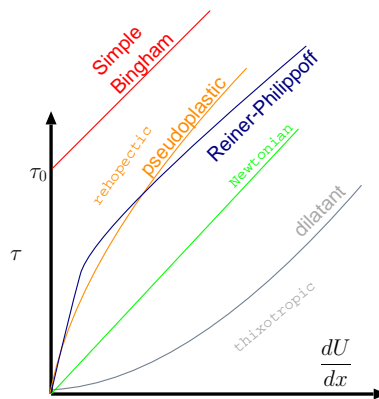


Fig. -1.5. The different of power fluids familys.

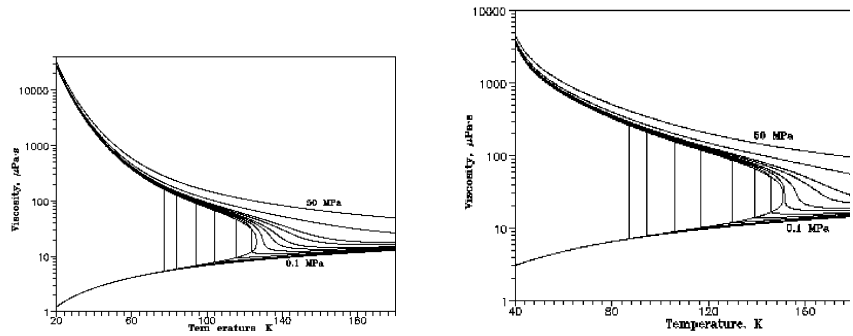


Fig. -1.6. Nitrogen (left) and Argon (right) viscosity as a function of the temperature and pressure after Lemmon and Jacobsen.

Figure 1.6 demonstrates that viscosity increases slightly with pressure, but this variation is negligible for most engineering problems. Well above the critical point, both materials are only a function of the temperature. On the liquid side below the critical point, the pressure has minor effect on the viscosity. It must be stress that the viscosity in the dome is meaningless. There is no such a thing of viscosity at 30% liquid. It simply depends on the structure of the flow as will be discussed in the chapter on multi phase flow. The lines in the above diagrams are only to show constant pressure lines. Oils have the greatest increase of viscosity with pressure which is a good thing for many engineering purposes.

1.5.2 Non-Newtonian Fluids

In equation (1.5), the relationship between the velocity and the shear stress was assumed to be linear. Not all the materials obey this relationship. There is a large class of materials which shows a non-linear relationship with velocity for any shear stress. This class of materials can be approximated by a single polynomial term that is $a = bx^n$. From the physical point of view, the coefficient depends on the velocity gradient. This relationship is referred to as power relationship and it can be written as

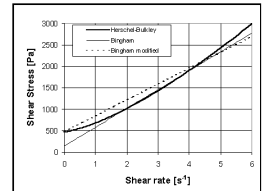


Fig. -1.7. The shear stress as a function of the shear rate.

$$\tau = K \overbrace{\left(\frac{dU}{dx}\right)^{n-1}}^{\text{viscosity}} \left(\frac{dU}{dx}\right) \quad (1.13)$$

The new coefficients (n, K) in equation (1.13) are constant. When $n = 1$ equation represent Newtonian fluid and K becomes the familiar μ . The viscosity coefficient is

always positive. When n , is above one, the liquid is dilatant. When n is below one, the fluid is pseudoplastic. The liquids which satisfy equation (1.13) are referred to as purely viscous fluids. Many fluids satisfy the above equation. Fluids that show increase in the viscosity (with increase of the shear) referred to as thixotropic and those that show decrease are called reoplectic fluids (see Figure 1.5).

Materials which behave up to a certain shear stress as a solid and above it as a liquid are referred to as Bingham liquids. In the simple case, the “liquid side” is like Newtonian fluid for large shear stress. The general relationship for simple Bingham flow is

$$\tau_{xy} = -\mu \pm \tau_0 \quad \text{if } |\tau_{yx}| > \tau_0 \quad (1.14)$$

$$\frac{dU_x}{dy} = 0 \quad \text{if } |\tau_{yx}| < \tau_0 \quad (1.15)$$

There are materials that simple Bingham model does not provide adequate explanation and a more sophisticated model is required. The Newtonian part of the model has to be replaced by power liquid. For example, according to Ferraris et al¹ concrete behaves as shown in Figure 1.7. However, for most practical purposes, this kind of figures isn't used in regular engineering practice.

1.5.3 Kinematic Viscosity

The kinematic viscosity is another way to look at the viscosity. The reason for this new definition is that some experimental data are given in this form. These results also explained better using the new definition. The kinematic viscosity embraces both the viscosity and density properties of a fluid. The above equation shows that the dimensions of ν to be square meter per second, $[m^2/sec]$, which are acceleration units (a combination of kinematic terms). This fact explains the name “kinematic” viscosity. The kinematic viscosity is defined as

$$\nu = \frac{\mu}{\rho} \quad (1.16)$$

The gas density increases with the temperature. The increase of the absolute viscosity with the temperature is enough to overcome the increase of density and thus, the kinematic viscosity also increase with the temperature for many materials.

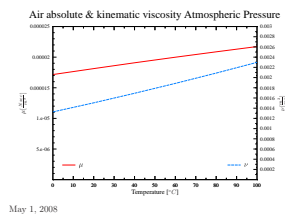


Fig. -1.8. Air viscosity as a function of the temperature.

¹C. Ferraris, F. de Larrard and N. Martys, Materials Science of Concrete VI, S. Mindess and J. Skalny, eds., 215-241 (2001)

1.5.4 Estimation of The Viscosity

The absolute viscosity of many fluids relatively doesn't change with the pressure but very sensitive to temperature. For isothermal flow, the viscosity can be considered constant in many cases. The variations of air and water as a function of the temperature at atmospheric pressure are plotted in Figures 1.8 and 1.9.

Some common materials (pure and mixture) have expressions that provide an estimate. For many gases, Sutherland's equation is used and according to the literature, provides reasonable results² for the range of $-40^{\circ}C$ to $1600^{\circ}C$

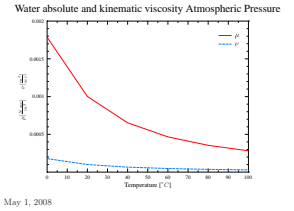


Fig. -1.9. Water viscosity as a function temperature.

$$\mu = \mu_0 \frac{0.555 T_{i0} + Suth}{0.555 T_{in} + Suth} \left(\frac{T}{T_0} \right)^{\frac{3}{2}} \quad (1.17)$$

Where

μ	viscosity at input temperature T
μ_0	reference viscosity at reference temperature, T_{i0}
T_{in}	input temperature in degrees Kelvin
T_{i0}	reference temperature in degrees Kelvin
$Suth$	Suth is Sutherland's constant and it is presented in the Table 1.1

Example 1.3:

Calculate the viscosity of air at 800K based on Sutherland's equation. Use the data provide in Table 1.1.

SOLUTION

Applying the constants from Sutherland's table provides

$$\mu = 0.00001827 \times \frac{0.555 \times 524.07 + 120}{0.555 \times 800 + 120} \times \left(\frac{800}{524.07} \right)^{\frac{3}{2}} \sim 2.51 \cdot 10^{-5} \left[\frac{N \cdot sec}{m^2} \right]$$

The viscosity increases almost by 40%. The observed viscosity is about $\sim 3.710^{-5}$.

End Solution

Liquid Metals

²This author is ambivalent about statement.

Material	coefficients	chemical formula	Sutherland	$T_{i0}[K]$	$\mu_0(N\ sec/m^2)$
ammonia		NH_3	370	527.67	0.00000982
standard air			120	524.07	0.00001827
carbon dioxide		CO_2	240	527.67	0.00001480
carbon monoxide		CO	118	518.67	0.00001720
hydrogen		H_2	72	528.93	0.0000876
nitrogen		N_2	111	540.99	0.0001781
oxygen		O_2	127	526.05	0.0002018
sulfur dioxide		SO_2	416	528.57	0.0001254

Table -1.1. The list for Sutherland's equation coefficients for selected materials.

Substance	formula	Temperature $T [^{\circ}C]$	Viscosity [$\frac{N\ sec}{m^2}$]
	$i - C_4 H_{10}$	23	0.0000076
	CH_4	20	0.0000109
	CO_2	20	0.0000146
oxygen	O_2	20	0.0000203
mercury vapor	Hg	380	0.0000654

Table -1.2. Viscosity of selected gases.

Substance	formula	Temperature T [$^{\circ}C$]	Viscosity [$\frac{N \cdot sec}{m^2}$]
	$(C_2H_5)O$	20	0.000245
	C_6H_6	20	0.000647
	Br_2	26	0.000946
	C_2H_5OH	20	0.001194
	Hg	25	0.001547
	H_2SO_4	25	0.01915
Olive Oil		25	0.084
Castor Oil		25	0.986
Clucose		25	5-20
Corn Oil		20	0.072
SAE 30		-	0.15-0.200
SAE 50		$\sim 25^{\circ}C$	0.54
SAE 70		$\sim 25^{\circ}C$	1.6
Ketchup		$\sim 20^{\circ}C$	0,05
Ketchup		$\sim 25^{\circ}C$	0,098
Benzene		$\sim 20^{\circ}C$	0.000652
Firm glass		-	$\sim 1 \times 10^7$
Glycerol		20	1.069

Table -1.3. Viscosity of selected liquids.

viscosity is obtained as

$$\mu_c = \frac{\overbrace{\mu}^{\text{given}}}{\underbrace{\mu_r}_{\text{figure 1.11}}} \quad (1.18)$$

The third way, when none is available, is by utilizing the following approximation

$$\mu_c = \sqrt{MT_c} \tilde{v}_c^{2/3} \quad (1.19)$$

Where \tilde{v}_c is the critical molecular volume and M is molecular weight. Or

$$\mu_c = \sqrt{MP_c} T_c^{2/3} T_c^{-1/6} \quad (1.20)$$

Calculate the reduced pressure and the reduced temperature and from the Figure 1.11 obtain the reduced viscosity.

Example 1.4:

Estimate the viscosity of oxygen, O_2 at $100^\circ C$ and $20[Bar]$.

SOLUTION

The critical condition of oxygen are $P_c = 50.35[Bar]$ $T_c = 154.4$ $\mu_c = 18 \left[\frac{N \text{ sec}}{m^2} \right]$ The value of the reduced temperature is

$$T_r \sim \frac{373.15}{154.4} \sim 2.41$$

The value of the reduced pressure is

$$P_r \sim \frac{20}{50.35} \sim 0.4$$

From Figure 1.11 it can be obtained $\mu_r \sim 1.2$ and the predicted viscosity is

$$\mu = \mu_c \overbrace{\left(\frac{\mu}{\mu_c} \right)}^{\text{Table}} = 18 \times 1.2 = 21.6[N \text{ sec}/m^2]$$

The observed value is $24[N \text{ sec}/m^2]$ ³.

End Solution

Viscosity of Mixtures

In general the viscosity of liquid mixture has to be evaluated experimentally. Even for homogeneous mixture, there isn't silver bullet to estimate the viscosity. In this book, only the mixture of low density gases is discussed for analytical expression. For most

³Kyama, Makita, Rev. Physical Chemistry Japan Vol. 26 No. 2 1956.

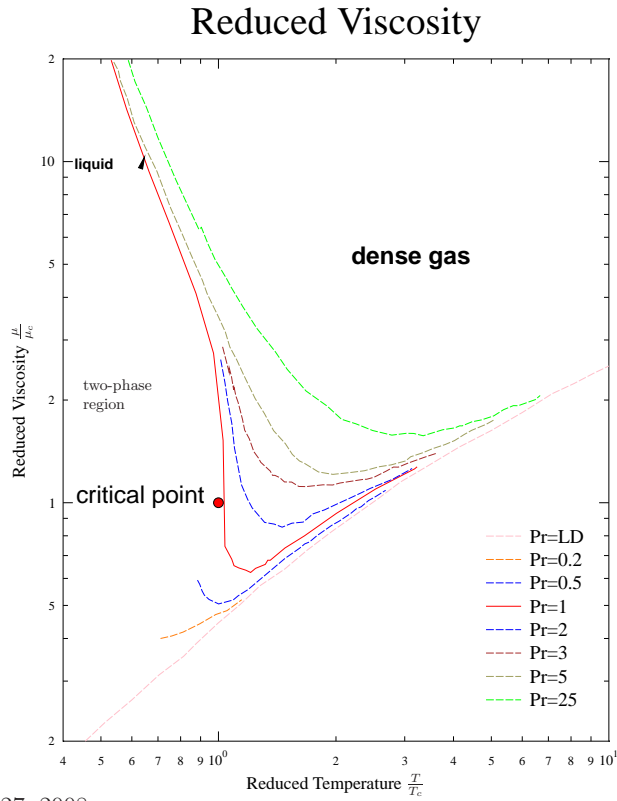


Fig. -1.11. Reduced viscosity as function of the reduced temperature.

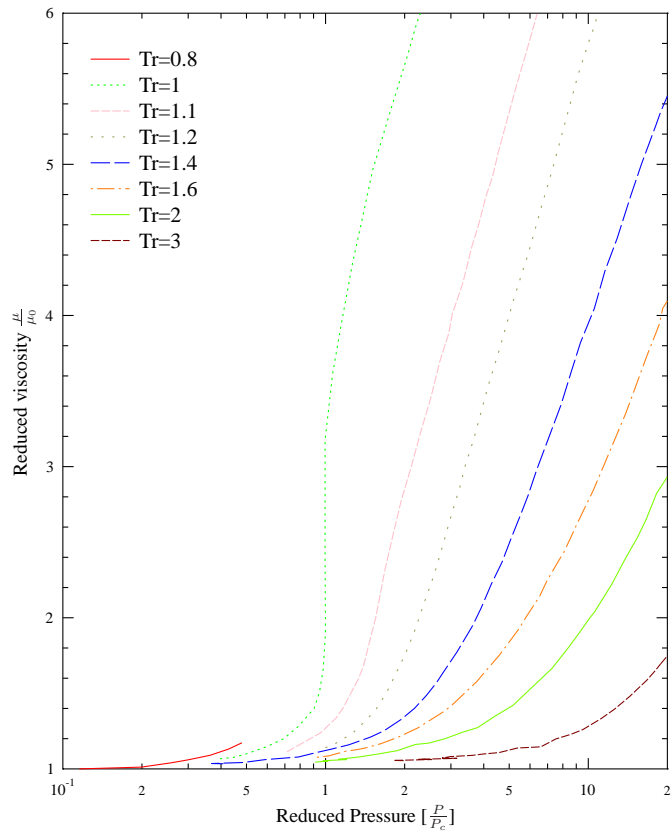
cases, the following Wilke's correlation for gas at low density provides a result in a reasonable range.

$$\mu_{mix} = \frac{\sum_{i=1}^n x_i \mu_i}{\sum_{j=1}^n x_j \Phi_{ij}} \quad (1.21)$$

where Φ_{ij} is defined as

$$\Phi_{ij} = \frac{1}{\sqrt{8}} \sqrt{1 + \frac{M_i}{M_j}} \left(1 + \sqrt{\frac{\mu_i}{\mu_j}} \sqrt{\frac{M_j}{M_i}} \right)^2 \quad (1.22)$$

Here, n is the number of the chemical components in the mixture. x_i is the mole fraction of component i , and μ_i is the viscosity of component i . The subscript i should be used for the j index. The dimensionless parameter Φ_{ij} is equal to one when $i = j$. The mixture viscosity is highly nonlinear function of the fractions of the components.



June 2, 2008

Fig. -1.12. Reduced viscosity as function of the reduced temperature.

Example 1.5:

Calculate the viscosity of a mixture (air) made of 20% oxygen, O_2 and 80% nitrogen N_2 for the temperature of $20^\circ C$.

SOLUTION

The following table summarizes the known details

i	Component	Molecular Weight, M	Mole Fraction, x	Viscosity, μ
1	O_2	32.	0.2	0.0000203
2	N_2	28.	0.8	0.00001754

i	j	M_i/M_j	μ_i/μ_j	Φ_{ij}
1	1	1.0	1.0	1.0
	2	1.143	1.157	1.0024
2	1	0.875	.86	0.996
	2	1.0	1.0	1.0

$$\mu_{mix} \sim \frac{0.2 \times 0.0000203}{0.2 \times 1.0 + 0.8 \times 1.0024} + \frac{0.8 \times 0.00001754}{0.2 \times 0.996 + 0.8 \times 1.0} \sim 0.0000181 \left[\frac{N \text{ sec}}{m^2} \right]$$

The observed value is $\sim 0.0000182 \left[\frac{N \text{ sec}}{m^2} \right]$.

End Solution

In very low pressure, in theory, the viscosity is only a function of the temperature with a "simple" molecular structure. For gases with very long molecular structure or complexity structure these formulas cannot be applied. For some mixtures of two liquids it was observed that at a low shear stress, the viscosity is dominated by a liquid with high viscosity and at high shear stress to be dominated by a liquid with the low viscosity liquid. The higher viscosity is more dominate at low shear stress. Reiner and Phillippoff suggested the following formula

$$\frac{dU_x}{dy} = \left(\frac{1}{\mu_\infty + \frac{\mu_0 - \mu_\infty}{1 + \left(\frac{\tau_{xy}}{\tau_s} \right)^2}} \right) \tau_{xy} \quad (1.23)$$

Where the term μ_∞ is the experimental value at high shear stress. The term μ_0 is the experimental viscosity at shear stress approaching zero. The term τ_s is the characteristic shear stress of the mixture. An example for values for this formula, for Molten Sulfur at temperature $120^\circ C$ are $\mu_\infty = 0.0215 \left(\frac{N \text{ sec}}{m^2} \right)$, $\mu_0 = 0.00105 \left(\frac{N \text{ sec}}{m^2} \right)$, and $\tau_s = 0.0000073 \left(\frac{kN}{m^2} \right)$. This equation (1.23) provides reasonable value only up to $\tau = 0.001 \left(\frac{kN}{m^2} \right)$.

Figure 1.12 can be used for a crude estimate of dense gases mixture. To estimate the viscosity of the mixture with n component Hougen and Watson's method for pseudocritical properties is adapted. In this method the following is defined as

$$P_{c_{mix}} = \sum_{i=1}^n x_i P_{c_i} \quad (1.24)$$

$$T_{c_{mix}} = \sum_{i=1}^n x_i T_{c_i} \quad (1.25)$$

and

$$\mu_{c, mix} = \sum_{i=1}^n x_i \mu_{c, i} \quad (1.26)$$

1.5.5 Bulk Modulus

Similar to solids (hook's law), liquids have a property that describes the volume change as results of pressure change for constant temperature. It can be noted that this property is not the result of the equation of state but related to it. The bulk modulus is defined as

$$B_T = -v \left(\frac{\partial P}{\partial v} \right)_T \quad (1.27)$$

Using the identity of $v = 1/\rho$ transfers equation (1.27) into

$$B_T = \rho \left(\frac{\partial P}{\partial \rho} \right)_T \quad (1.28)$$

The bulk modulus for several liquids is presented in Table 1.5.

Table -1.5. The bulk modulus for selected material with the critical temperature and pressure *na* → not available and *nf* → not found (exist but was not found in the literature).

chemical component	Bulk Modulus $10^9 \frac{N}{m}$	T_c	P_c
Acetic Acid	2.49	593K	57.8 [Bar]
Acetone	0.80	508 K	48 [Bar]
Benzene	1.10	562 K	4.74 [MPa]
Carbon Tetrachloride	1.32	556.4 K	4.49 [MPa]
Ethyl Alcohol	1.06	514 K	6.3 [Mpa]
Gasoline	1.3	nf	nf
Glycerol	4.03-4.52	850 K	7.5 [Bar]
Mercury	26.2-28.5	1750 K	172.00 [MPa]
Methyl Alcohol	0.97	Est 513	Est 78.5 [Bar]
Nitrobenzene	2.20	nf	nf
Olive Oil	1.60	nf	nf
Paraffin Oil	1.62	nf	nf
SAE 30 Oil	1.5	na	na
Seawater	2.34	na	na
Toluene	1.09	591.79 K	4.109 [MPa]
Turpentine	1.28	na	na
Water	2.15-2.174	647.096 K	22.064 [MPa]

In the literature, additional expansions for similar parameters are defined. The thermal expansion is defined as

$$\beta_P = \frac{1}{v} \left(\frac{\partial v}{\partial T} \right)_P \tag{1.29}$$

This parameter indicates the change of volume due to temperature change when the pressure is constant. Another definition is referred as coefficient of tension and it is defined as

$$\beta_v = \frac{1}{P} \left(\frac{\partial P}{\partial T} \right)_v \tag{1.30}$$

This parameter indicates the change of the pressure due to the change of temperature (where $v = \text{constant}$). These definitions are related to each other. This relationship is obtained by the observation that the pressure as a function of the temperature and specific volume as

$$P = f(T, v) \tag{1.31}$$

The full pressure derivative is

$$dP = \left(\frac{\partial P}{\partial T} \right)_v dT + \left(\frac{\partial P}{\partial v} \right)_T dv \tag{1.32}$$

On constant pressure lines, $dP = 0$, and therefore equation (1.32) is

$$0 = \left(\frac{\partial P}{\partial T} \right)_v dT + \left(\frac{\partial P}{\partial v} \right)_T dv \tag{1.33}$$

From equation (1.33) follows that

$$\left. \frac{dv}{dT} \right|_{P=\text{const}} = - \frac{\left(\frac{\partial P}{\partial T} \right)_v}{\left(\frac{\partial P}{\partial v} \right)_T} \tag{1.34}$$

Equation (1.34) indicates that relationship for these three coefficients is

$$\beta_T = - \frac{\beta_v}{\beta_P} \tag{1.35}$$

The last equation (1.35) sometimes is used in measurement of the bulk modulus.

The increase of the pressure increases the bulk modulus due to the molecules increase of the rejecting forces between each other when they are closer. In contrast,

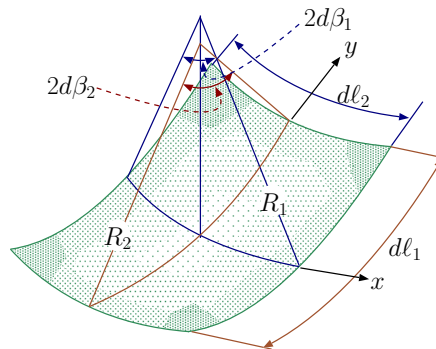


Fig. -1.13. Surface Tension control volume analysis.

the temperature increase results in reduction of the bulk modulus because the molecules are further away.

Example 1.6:

Calculate the modulus of liquid elasticity that reduced 0.035 per cent of its volume by applying a pressure of 5[Bar].

SOLUTION

Using the definition for the bulk modulus

$$\beta_T = -v \frac{\partial P}{\partial v} \sim \frac{v}{\Delta v} \Delta P = \frac{5}{0.00035} \sim 14285.714[\text{Bar}]$$

End Solution

Example 1.7:

Calculate the pressure needed to apply on water to reduce its volume by 1 per cent. Assume the temperature to be 20°C.

SOLUTION

Using the definition for the bulk modulus

$$\Delta P \sim \beta_T \frac{\Delta v}{v} \sim 2.15 \cdot 10^9 \cdot 0.01 = 2.15 \cdot 10^7 [\text{N/m}^2] = 215[\text{Bar}]$$

End Solution

1.6 Surface Tension

The surface tension manifested itself by a rise or depression of the liquid at the free surface edge. Surface tension is also responsible for the creation of the drops and bubbles. It is also responsible for the breakage of a liquid jet into other medium/phase to many drops (atomization). The surface tension is force per length and is measured by [N/m] and is acting to stretch the surface.

Surface tension results from a sharp change in the density between two adjacent phases or materials. There is a common misconception for the source of the surface tension. In many (physics, surface tension, and fluid mechanics) books explained that the surface tension is a result from unbalanced molecular cohesive forces. This explanation is wrong since it is in conflict with Newton's second law (see example ?). This erroneous explanation can be traced to Adam's book but earlier source may be found.

The relationship between the surface tension and the pressure on the two sides of the surface is based on geometry. Consider a small element of surface. The pressure

on one side is P_i and the pressure on the other side is P_o . When the surface tension is constant, the horizontal forces cancel each other because symmetry. In the vertical direction, the surface tension forces are pulling the surface upward. Thus, the pressure difference has to balance the surface tension. The forces in the vertical direction reads

$$(P_i - P_o) dl_1 dl_2 = \Delta P dl_1 dl_2 = 2\sigma dl_1 \sin \beta_1 + 2\sigma dl_2 \sin \beta_2 \quad (1.36)$$

For a very small area, the angles are very small and thus $(\sin \beta \sim \beta)$. Furthermore, it can be noticed that $dl_i \sim 2 R_i d\beta_i$. Thus, the equation (1.36) can be simplified as

$$\Delta P = \sigma \left(\frac{1}{R_1} + \frac{1}{R_2} \right) \quad (1.37)$$

Equation (1.37) predicts that pressure difference increase with inverse of the radius. There are two extreme cases: one) radius of infinite and radius of finite size. The second with two equal radii. The first case is for an infinite long cylinder for which the equation (1.37) is reduced to

$$\Delta P = \sigma \left(\frac{1}{R} \right) \quad (1.38)$$

Other extreme is for a sphere for which the main radii are the same and equation (1.37) is reduced to

$$\Delta P = \frac{2\sigma}{R} \quad (1.39)$$

Where R is the radius of the sphere. A soap bubble is made of two layers, inner and outer, thus the pressure inside the bubble is

$$\Delta P = \frac{4\sigma}{R} \quad (1.40)$$

Example 1.8:

A Tank filled with liquid, which contains n bubbles with equal radii, r . Calculate the minimum work required to increase the pressure in tank by ΔP . Assume that the liquid bulk modulus is infinity.

SOLUTION

The work is due to the change of the bubbles volume. The work is

$$w = \int_{r_0}^{r_f} \Delta P(v) dv \quad (1.41)$$

The minimum work will be for a reversible process. The reversible process requires very slow compression. It is worth noting that for very slow process, the temperature must

remain constant due to heat transfer. The relationship between pressure difference and the radius is described by equation (1.39) for reversible process. Hence the work is

$$w = \int_{r_0}^{r_f} \underbrace{\frac{\Delta P}{r}}_{\frac{2\sigma}{r}} \underbrace{4\pi r^2 dr}_{dv} = 8\pi\sigma \int_{r_0}^{r_f} r dr = 4\pi\sigma (r_f^2 - r_0^2) \quad (1.42)$$

Where, r_0 is the radius at the initial stage and r_f is the radius at the final stage.

The work for n bubbles is then $4\pi\sigma n (r_f^2 - r_0^2)$. It can be noticed that the work is negative, that is the work is done on the system.

End Solution

1.6.1 Wetting of Surfaces

To explain the source of the contact angle, consider the point where three phases became in contact. This contact point occurs due to free surface reaching a solid boundary. The surface tension occurs between gas phase (G) to liquid phase (L) and also occurs between the solid (S) and the liquid phases as well as between the gas phase and the solid phase. In Figure 1.14, forces diagram is shown when control volume is chosen so that the masses of the solid, liquid, and gas can be ignored. Regardless to the magnitude of the surface tensions (except to zero) the forces cannot be balanced for the description of straight lines. For example, forces balanced along the line of solid boundary is

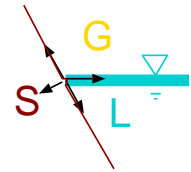


Fig. -1.14. Forces in Contact angle.

$$\sigma_{gs} - \sigma_{ls} - \sigma_{lg} \cos \beta = 0 \quad (1.43)$$

and in the tangent direction to the solid line the forces balance is

$$F_{solid} = \sigma_{lg} \sin \beta \quad (1.44)$$

substituting equation (1.44) into equation (1.43) yields

$$\sigma_{gs} - \sigma_{ls} = \frac{F_{solid}}{\tan \beta} \quad (1.45)$$

For $\beta = \pi/2 \implies \tan \beta = \infty$. Thus, the solid reaction force must be zero. The gas solid surface tension is different from the liquid solid surface tension and hence violating equation (1.43).

The surface tension forces must be balanced, thus, a contact angle is created to balance it. The contact angle is determined by whether the surface tension between the gas

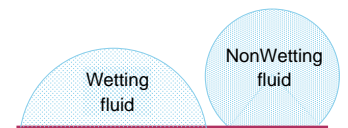


Fig. -1.15. Description of wetting and non-wetting fluids.

solid (gs) is larger or smaller than the surface tension of liquid solid (ls) and the local geometry. It must be noted that the solid boundary isn't straight. The surface tension is a molecular phenomenon, thus depends on the local structure of the surface and it provides the balance for these local structures.

The connection of the three phases-materials-mediums creates two situations which are categorized as wetting or non-wetting. There is a common definition of wetting the surface. If the angle of the contact between three materials is larger than 90° then it is non-wetting. On the other hand, if the angle is below 90° the material is wetting the surface (see Figure 1.15). The angle is determined by properties of the liquid, gas medium and the solid surface. And a small change on the solid surface can change the wetting condition to non-wetting. In fact there are commercial sprays that are intent to change the surface from wetting to non wetting. This fact is the reason that no reliable data can be provided with the exception to pure substances and perfect geometries. For example, water is described in many books as a wetting fluid. This statement is correct in most cases, however, when solid surface is made or coated with certain materials, the water is changed to be non-wetting (for example 3M selling product to "change" water to non-wetting). So, the wetness of fluids is a function of the solid as well.

Table -1.6. The contact angle for air, distilled water with selected materials to demonstrate the inconsistency.

chemical component	Contact Angle	Source
Steel	$\pi/3.7$	[1]
Steel,Nickel	$\pi/4.74$	[2]
Nickel	$\pi/4.74$ to $\pi/3.83$	[1]
Nickel	$\pi/4.76$ to $\pi/3.83$	[3]
Chrome-Nickel Steel	$\pi/3.7$	[4]
Silver	$\pi/6$ to $\pi/4.5$	[5]
Zink	$\pi/3.4$	[4]
Bronze	$\pi/3.2$	[4]
Copper	$\pi/4$	[4]
Copper	$\pi/3$	[7]
Copper	$\pi/2$	[8]

- 1 R. Siegel, E. G. Keshock (1975) "Effects of reduced gravity on nucleate boiling bubble dynamics in saturated water," AIChE Journal Volume 10 Issue 4, Pages 509 - 517. 1975
- 2 Bergles A. E. and Rohsenow W. M. "The determination of forced convection surface-

- boiling heat transfer, ASME, J. Heat Transfer, vol 1 pp 365 - 372.
- 3 Tolubinsky, V.I. and Ostrovsky, Y.N. (1966) "On the mechanism of boiling heat transfer", International Journal of Heat and Mass Transfer, Vol. 9, No 12, pages 1465-1470.
 - 4 Arefeva E.I., Aladev O, I.T., (1958) "wlijanii smatchivaemosti na teploobmen pri kipenii," Injenerno Fizitcheskij Jurnal, 11-17 1(7) In Russian.
 - 5 Labuntsov D. A. (1939) "Approximate theory of heat transfer by developed nucleate boiling" In Sussian Izvestiya An SSSR , Energetika I transport, No 1.
 - 6 Basu, N., Warriar, G. R., and Dhir, V. K., (2002) "Onset of Nucleate Boiling and Active Nucleation Site Density during Subcooled Flow Boiling," ASME Journal of Heat Transfer, Vol. 124, pages 717 -728.
 - 7 Gaetner, R. F., and Westwater, J. W., (1960) "Population of Active Sites in Nucleate Boiling Heat Transfer," Chem. Eng. Prog. Symp., Ser. 56.
 - 8 Wang, C. H., and Dhir, V. K., (1993), "Effect of Surface Wettability on Active Nucleation Site Density During Pool Boiling of Water on a Vertical Surface," J. Heat Transfer 115, pp. 659-669

To explain the contour of the surface, and the contact angle consider simple "wetting" liquid contacting a solid material in two-dimensional shape as depicted in Figure 1.16. To solve the shape of the liquid surface, the pressure difference between the two sides of free surface has to be balanced by the surface tension. In Figure 1.16 describes the raising of the liquid as results of the surface tension. The surface tension reduces the pressure in the liquid above the liquid line (the dotted line in the Figure 1.16). The pressure just below the surface is $-g h(x) \rho$ (this pressure difference will be explained in more details in Chapter 4). The pressure, on the gas side, is the atmospheric pressure. This problem is a two dimensional problem and equation (1.38) is applicable to it. Applying equation (1.38) and using the pressure difference yields

$$g h(x), \rho = \frac{\sigma}{R(x)} \quad (1.46)$$

The radius of any continuous function, $h = h(x)$, is

$$R(x) = \frac{\left(1 + [\dot{h}(x)]^2\right)^{3/2}}{\ddot{h}(x)} \quad (1.47)$$

Where \dot{h} is the derivative of h with respect to x .

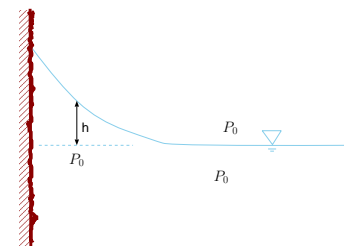


Fig. -1.16. Description of liquid surface.

Equation (1.47) can be derived either by forcing a circle at three points at $(x, x+dx, \text{ and } x+2dx)$ and thus finding the the diameter or by geometrical analysis of triangles build on points x and $x+dx$ (perpendicular to the tangent at these points). Substituting equation (1.47) into equation (1.46) yields

$$g h(x) \rho = \frac{\sigma}{\frac{\left(1 + \left[\dot{h}(x)\right]^2\right)^{3/2}}{\ddot{h}(x)}} \quad (1.48)$$

Equation (1.48) is non-linear differential equation for height and can be written as

$$\frac{g h \rho}{\sigma} \left(1 + \left[\frac{dh}{dx}\right]^2\right)^{3/2} - \frac{d^2 h}{dx^2} = 0 \quad (1.49)$$

With the boundary conditions that specify either the derivative $\dot{h}(x = r) = 0$ (symmetry) and the derivative at $\dot{h}x = \beta$ or heights in two points or other combinations. An alternative presentation of equation (1.48) is

$$g h \rho = \frac{\sigma \ddot{h}}{\left(1 + \dot{h}^2\right)^{3/2}} \quad (1.50)$$

Integrating equation (1.50) transforms into

$$\int \frac{g \rho}{\sigma} h dh = \int \frac{\ddot{h}}{\left(1 + \dot{h}^2\right)^{3/2}} dh \quad (1.51)$$

The constant $Lp \sigma / \rho g$ is referred to as Laplace's capillarity constant. The units of this constant are meter squared. The differential dh is \dot{h} . Using dummy variable and the identys $\dot{h} = \xi$ and hence, $\ddot{h} = \dot{\xi} = d\xi$ transforms equation (1.51) into

$$\int \frac{1}{Lp} h dh = \int \frac{\xi d\xi}{\left(1 + \xi^2\right)^{3/2}} \quad (1.52)$$

After the integration equation (1.52) becomes

$$\frac{h^2}{2 Lp} + constant = -\frac{1}{\left(1 + \dot{h}^2\right)^{1/2}} \quad (1.53)$$

At infinity, the height and the derivative of the height must be zero so $constant + 0 = -1/1$ and hence, $constant = -1$.

$$1 - \frac{h^2}{2Lp} = \frac{1}{(1 + \dot{h}^2)^{1/2}} \quad (1.54)$$

Equation (1.54) is a first order differential equation that can be solved by variables separation⁴. Equation (1.54) can be rearranged to be

$$(1 + \dot{h}^2)^{1/2} = \frac{1}{1 - \frac{h^2}{2Lp}} \quad (1.55)$$

Squaring both sides and moving the one to the right side yields

$$\dot{h}^2 = \left(\frac{1}{1 - \frac{h^2}{2Lp}} \right)^2 - 1 \quad (1.56)$$

The last stage of the separation is taking the square root of both sides to be

$$\dot{h} = \frac{dh}{dx} = \sqrt{\left(\frac{1}{1 - \frac{h^2}{2Lp}} \right)^2 - 1} \quad (1.57)$$

or

$$\frac{dh}{\sqrt{\left(\frac{1}{1 - \frac{h^2}{2Lp}} \right)^2 - 1}} = dx \quad (1.58)$$

Equation (1.58) can be integrated to yield

$$\int \frac{dh}{\sqrt{\left(\frac{1}{1 - \frac{h^2}{2Lp}} \right)^2 - 1}} = x + constant \quad (1.59)$$

The constant is determined by the boundary condition at $x = 0$. For example if $h(x = 0) = h_0$ then $constant = h_0$. This equation is studied extensively in classes on surface tension. Furthermore, this equation describes the dimensionless parameter that affects this phenomenon and this parameter will be studied in Chapter ?. This book is introductory, therefore this discussion on surface tension equation will be limited.

1.6.1.1 Capillarity

⁴This equation has an analytical solution which is $x = Lp\sqrt{4 - (h/Lp)^2} - Lp \operatorname{acosh}(2Lp/h) + constant$ where Lp is the Laplace constant. Shamefully, this author doesn't know how to show it in a two lines derivations.

1.6. SURFACE TENSION

The capillary forces referred to the fact that surface tension causes liquid to rise or penetrate into area (volume), otherwise it will not be there. It can be shown that the height that the liquid raised in a tube due to the surface tension is

$$h = \frac{2\sigma \cos \beta}{g \Delta \rho r} \quad (1.60)$$

Where $\Delta \rho$ is the difference of liquid density to the gas density and r is the radius of tube.

But this simplistic equation is unusable and useless unless the contact angle (assuming that the contact angle is constant or a repressive average can be found or provided or can be measured) is given. However, in reality there is no readily information for contact angle⁵ and therefore this equation is useful to show the trends. The maximum that the contact angle can be obtained in equation (1.60) when $\beta = 0$ and thus $\cos \beta = 1$. This angle is obtained when a perfect half a sphere shape exist of the liquid surface. In that case equation (1.60) becomes

$$h_{max} = \frac{2\sigma}{g \Delta \rho r} \quad (1.61)$$

Figure 1.18 exhibits the height as a function of the radius of the tube. The height based on equation (1.61) is shown in Figure 1.17 as blue line. The actual height is shown in the red line. Equation (1.61) provides reasonable results only in a certain range. For a small tube radius, equation (1.49) proved better results because the curve approaches hemispherical shaper (small gravity effect). For large radii equation (1.49) approaches the strait line (the liquid line) strong gravity effect. On the other hand, for extremely small radii equation (1.61) indicates that the high height which indicates a negative pressure. The liquid at a certain pressure will be vaporized and will breakdown the model upon this equation was constructed. Furthermore, the small scale indicates that the simplistic and continuous approach is not appropriate and a different model is needed. The conclusion of this discussion are shown in Figure 1.17. The actual dimension for many liquids (even water) is about 1-5 [mm].

The discussion above was referred to "wetting" contact angle. The depression of the liquid occurs in a "negative" contact angle similarly to "wetting." The depression height, h is similar to equation (1.61) with a minus sign. However, the gravity is working against the surface tension and reducing the range and quality of the predictions of equation (1.61). The measurements of the height of distilled water and mercury are

⁵Actually, there are information about the contact angle. However, that information conflict each other and no real information is available see Table 1.6.

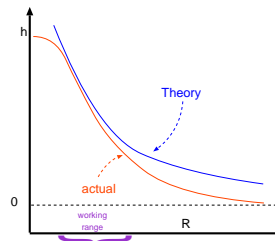
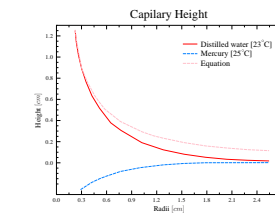


Fig. -1.17. The raising height as a function of the radii.



May 29, 2008

Fig. -1.18. The raising height as a function of the radius.

presented in Figure 1.18. The experimental results of these materials are with agreement with the discussion above.

The surface tension of a selected material is given in Table 1.7.

In conclusion, the surface tension issue is important only in case where the radius is very small and gravity is negligible. The surface tension depends on the two materials or mediums that it separates.

Example 1.9:

Calculate the diameter of a water droplet to attain pressure difference of $1000[N/m^2]$. You can assume that temperature is $20^\circ C$.

SOLUTION

The pressure inside the droplet is given by equation (1.39).

$$D = 2R = \frac{2\sigma}{\Delta P} = \frac{4 \times 0.0728}{1000} \sim 2.912 \cdot 10^{-4} [m]$$

End Solution

Example 1.10:

Calculate the pressure difference between a droplet of water at $20^\circ C$ when the droplet has a diameter of 0.02 cm .

SOLUTION

using equation

$$\Delta P = \frac{2\sigma}{r} \sim \frac{2 \times 0.0728}{0.0002} \sim 728.0 [N/m^2]$$

End Solution

Example 1.11:

Calculate the maximum force necessary to lift a thin wire ring of $0.04[m]$ diameter from a water surface at $20^\circ C$. Neglect the weight of the ring.

SOLUTION

$$F = 2(2\pi r \sigma) \cos \beta$$

The actual force is unknown since the contact angle is unknown. However, the maximum Force is obtained when $\beta = 0$ and thus $\cos \beta = 1$. Therefore,

$$F = 4\pi r \sigma = 4 \times \pi \times 0.04 \times 0.0728 \sim .0366 [N]$$

In this value the gravity is not accounted for.

End Solution

Table -1.7. The surface tension for selected materials at temperature 20°C when not mentioned.

chemical component	Surface Tension $\frac{mN}{m}$	T	correction $\frac{mN}{m K}$
Acetic Acid	27.6	20°C	n/a
Acetone	25.20	-	-0.1120
Aniline	43.4	22°C	-0.1085
Benzene	28.88	-	-0.1291
Benzylalcohol	39.00	-	-0.0920
Benzylbenzoate	45.95	-	-0.1066
Bromobenzene	36.50	-	-0.1160
Bromobenzene	36.50	-	-0.1160
Bromoform	41.50	-	-0.1308
Butyronitrile	28.10	-	-0.1037
Carbon disulfid	32.30	-	-0.1484
Quinoline	43.12	-	-0.1063
Chloro benzene	33.60	-	-0.1191
Chloroform	27.50	-	-0.1295
Cyclohexane	24.95	-	-0.1211
Cyclohexanol	34.40	25°C	-0.0966
Cyclopentanol	32.70	-	-0.1011
Carbon Tetrachloride	26.8	-	n/a
Carbon disulfid	32.30	-	-0.1484
Chlorobutane	23.10	-	-0.1117
Ethyl Alcohol	22.3	-	n/a
Ethanol	22.10	-	-0.0832
Ethylbenzene	29.20	-	-0.1094
Ethylbromide	24.20	-	-0.1159
Ethylene glycol	47.70	-	-0.0890
Formamide	58.20	-	-0.0842
Gasoline	~ 21	-	n/a
Glycerol	64.0	-	-0.0598
Helium	0.12	-269°C	n/a
Mercury	425-465.0	-	-0.2049
Methanol	22.70	-	-0.0773
Methyl naphthalene	38.60	-	-0.1118
Methyl Alcohol	22.6	-	n/a
Neon	5.15	-247°C	n/a
Nitrobenzene	43.90	-	-0.1177
Olive Oil	43.0-48.0	-	-0.067
Perfluoroheptane	12.85	-	-0.0972
Perfluorohexane	11.91	-	-0.0935

Table -1.7. The surface tension for selected materials (continue)

chemical component	Surface Tension $\frac{mN}{m}$	T	correction $\frac{mN}{m K}$
Perfluorooctane	14.00	-	-0.0902
Phenylisothiocyanate	41.50	-	-0.1172
Propanol	23.70	$25^{\circ}C$	-0.0777
Pyridine	38.00	-	-0.1372
Pyrrol	36.60	-	-0.1100
SAE 30 Oil	n/a	-	n/a
Seawater	54-69	-	n/a
Toluene	28.4	-	-0.1189
Turpentine	27	-	n/a
Water	72.80	-	-0.1514
o-Xylene	30.10	-	-0.1101
m-Xylene	28.90	-	-0.1104